

ON STRONG SOLVABILITY OF ONE NONLOCAL BOUNDARY VALUE PROBLEM FOR POISSON EQUATION IN RECTANGLE

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Abstract

One nonlocal boundary value problem for the Poisson equation in an rectangular domain is considered in this work. The concept of strong solution of this problem is introduced. The correct solvability of this problem in weighted Sobolev spaces generated by a mixed norm is proved using the Fourier method. This is accomplished by exploiting the basis property of the system of eigenfunctions and associated functions of the corresponding spectral problem in weighted Lebesgue spaces, as well as the Hausdorff-Young inequality valid for this system.

Keywords: Poisson equation; nonlocal problem; weighted Sobolev space; strong solution

1. Introduction

Consider the following (formal for now) nonlocal boundary value problem for

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the Poisson equation:

$$\left. \begin{aligned} u_{xx} + u_{yy} &= f(x; y), \quad 0 < x < 2\pi, \quad 0 < y < h, \\ u(x, 0) &= \varphi(x), \quad u(x, h) = \psi(x), \quad 0 < x < 2\pi, \\ u_x(0, y) &= 0, \quad u(0, y) = u(2\pi, y), \quad 0 < y < h, \end{aligned} \right\} \quad (2)$$

In the case of a semi-infinite strip for one class of degenerate elliptic equation, which, in particular, covers the Laplace equation, this problem was previously considered by E.I. Moiseev [1] in the classical formulation. A similar problem was considered by M.E. Lerner and O.A. Repin [2]. The case of a semi-infinite strip for the Laplace equation is also considered in [3], where the strong solvability of a nonlocal boundary value problem in a weighted grand-Sobolev space is studied.

It is obvious that the solution to problem (1)-(2) can be sought as the sum of the solutions of the following two problems:

$$\left. \begin{aligned} u_{xx} + u_{yy} &= 0 \\ u(x, 0) &= \varphi(x), \quad u(x, h) = \psi(x), \\ u_x(0, y) &= 0, \quad u(0, y) = u(2\pi, y); \end{aligned} \right\} \quad (A)$$

and

$$\left. \begin{aligned} u_{xx} + u_{yy} &= f(x, y), \\ u(x, 0) &= 0, \quad u(x, h) = 0, \\ u_x(0, y) &= 0, \quad u(0, y) = u(2\pi, y). \end{aligned} \right\} \quad (B)$$

The strong solvability of Problem (A) in the domain $\Pi = (0, 2\pi) \times (0, h)$ in weighted Sobolev spaces was considered in [4]. In this paper, we study the question of the strong solvability of Problem (B) in a weighted Sobolev space with a weight from the Muckenhoupt class, which, together with the result of [4], will allow us to formulate a well-posed solvability of Problem (1)–(2) in this space. The concept of a strong solution is defined, and the well-posed solvability of Problem (B) is proved by the spectral method.

Such problems have specific peculiarities compared to the ones with local conditions. Earlier, F.I. Frankl [5]; [6, p.453-456] considered the problem with nonlocal boundary condition for a mixed type equation. Bitsadze-Samarski problem [7] for elliptic equations is also nonlocal with supports on a part of the boundary of domain, and these supports are free of other boundary conditions. In [8], N.I. Ionkin and E.I. Moiseev solved the boundary value problem for multi-

dimensional parabolic equations with nonlocal conditions, whose supports are the characteristic and the improper parts of the boundary of domain.

We note that the study of the solvability of elliptic equations with respect to the so-called non-standard function spaces, as well as in weighted Sobolev spaces, encounters certain difficulties compared to the weightless case. Therefore, the number of studies devoted to this area has been growing in recent years (see, for example, [3, 9 – 24]), and the development of the corresponding theory is far from complete. These attempts are also relevant in connection with the question of the basis properties of eigenfunctions and associated functions of spectral problems obtained by applying the method of separation of variables in the corresponding non-standard function spaces. The question of the basis property of the classical exponential system in weighted Lebesgue spaces was studied in [25, 26]. In this regard, we note the works [27 - 33], in which similar questions are studied and appropriate research methods are proposed.

2. Auxiliary concepts and facts

2.1. Notations.

We will use standard notations. N will be the set of positive integers, while $\alpha = (\alpha_1; \alpha_2) \in Z_+ \times Z_+$ will denote a multi-index, where $Z_+ = N \cup \{0\}$. Denote $\partial^\alpha u = \frac{\partial^{|\alpha|} u}{\partial x^{\alpha_1} \partial y^{\alpha_2}}$, where $|\alpha| = \alpha_1 + \alpha_2$. By $|M|$ we will denote the Lebesgue measure of the set M ; \bar{M} will be the closure of M . $C^\infty(\bar{M})$ will stand for the infinitely differentiable functions on \bar{M} , and $C_0^\infty(M)$ will denote the infinitely differentiable and finite functions on M . Throughout this paper we will assume that p' is a conjugate number of p , $1 < p < +\infty$: $\frac{1}{p'} + \frac{1}{p} = 1$.

2.2. Weighted Sobolev space $W_{p,\nu}^2(\Pi)$.

Let's define our weighted Sobolev space. Let $\nu: [0, 2\pi] \rightarrow (0, +\infty)$ be some weight function, $\Pi = (0, 2\pi) \times (0, h)$. Denote by $L_{p,\nu}(\Pi)$ a Banach space of functions on Π with the mixed norm

$$\|f\|_{L_{p,\nu}(\Pi)} = \int_0^h \left(\int_0^{2\pi} |f(x; y)|^p \nu(x) dx \right)^{\frac{1}{p}} dy, \quad 1 < p < +\infty.$$

Denote by $W_{p,\nu}^2(\Pi)$ a Sobolev space generated by the norm

$$\|u\|_{W_{p,\nu}^2(\Pi)} = \sum_{|\alpha| \leq 2} \|\partial^\alpha u\|_{L_{p,\nu}(\Pi)}.$$

Now denote by $L_{p,\nu}(I)$, where $I = (0, 2\pi)$, a weighted Lebesgue space generated by the norm

$$\|f\|_{L_{p,\nu}(I)} = \left(\int_I |f(x)|^p \nu(x) dx \right)^{\frac{1}{p}}.$$

We will also consider the weighted Sobolev space $W_{p,\nu}^2(I)$, generated by the norm

$$\|f\|_{W_{p,\nu}^2(I)} = \|f\|_{L_{p,\nu}(I)} + \|f'\|_{L_{p,\nu}(I)} + \|f''\|_{L_{p,\nu}(I)}.$$

We will need the class of Muckenhoupt weights $A_p(I)$ (see [26]). This is a class of 2π – periodic functions (i.e. the class of functions ν periodically extended to the whole of the axis with period 2π), satisfying the condition

$$\sup_{J \subset I} \left(\frac{1}{|J|} \int_J \nu(t) dt \right) \left(\frac{1}{|J|} \int_I |\nu(t)|^{-\frac{1}{p-1}} \right)^{p-1} < +\infty,$$

where sup is taken over all intervals $J \subset I$, and $|J|$ is a length of the interval J .

2.3. Statement of the problem.

For correct statement of considered boundary value problem we need firstly define the trace operator corresponding to the weighted Sobolev space $W_{p,\nu}^2(\Pi)$ on domain $\Pi = (0, 2\pi) \times (0, h) \subset R^2$ with boundary $\partial\Pi = I_0 \cup I_h \cup \Gamma_0 \cup \Gamma_{2\pi}$, where are denoted $I_0 = \{(x; 0) : 0 < x < 2\pi\}$, $I_h = \{(x; h) : 0 < x < 2\pi\}$, $\Gamma_0 = \{(0; y) : 0 < y < h\}$ and $\Gamma_{2\pi} = \{(2\pi; y) : 0 < y < h\}$. It is obvious that $L_{p,\nu}(I) \subset L_1(I)$. Then, the continuous embedding $W_{p,\nu}^2(\Pi) \subset W_1^2(\Pi)$ is also true. Consequently, every function $u \in W_{p,\nu}^2(\Pi)$ has traces $u|_{\partial\Pi}$ and $u_x|_{\partial\Pi}$ as functions of space $L_1(\partial\Pi; d\sigma)$ on the boundary (correctly defined with respect to the Lebesgue measure on $\partial\Pi$).

Consider the following nonlocal problem:

$$\Delta u = f(x; y), \quad (x; y) \in \Pi, \quad (3)$$

$$u|_{I_0} = 0, \quad u|_{I_h} = 0, \quad u|_{\Gamma_0} = u|_{\Gamma_{2\pi}}, \quad u_x|_{\Gamma_0} = 0, \quad (4)$$

where $f(x; y) \in L_{p,v}(\Pi)$. By the solution of this problem, we mean a function $u \in W_{p,v}^2(\Pi)$, which satisfies the equality (3) a.e. in Π and whose traces satisfy the relations (4) on the boundary $\partial\Pi$. Taking into account the above, we arrive at the following definition of strong solvability of problem (3)-(4).

Definition 2.1. A function $u \in W_{p,v}^2(\Pi)$ is called a strong solution of the problem (3)-(4) if the equality (3) is satisfied for a.e. $(x; y) \in \Pi$ and its trace $u|_{\partial\Pi}$ satisfies the relations (4).

2.4. Some facts from basis theory.

Let us present some well-known concepts and facts from basis theory. Let X be a Banach space. A system $\{u_n\}_{n \in \mathbb{N}} \subset X$ is called a basis if any element $f \in X$ is uniquely represented as a series

$$f = \sum_{n=1}^{\infty} c_n u_n$$

convergent in the norm X . A system $\{u_n\}_{n \in \mathbb{N}} \subset X$ is called complete in X if $\overline{\text{Sp}}\{u_n\} = X$ and minimal in X if $u_n \notin \overline{\text{Sp}}\{u_k\}_{k \neq n}$. It is known that each basis of the space X is a complete and minimal system in X , the converse is not true in general. The system $\{x_n\}_{n \in \mathbb{N}} \subset X$ is complete in $X \Leftrightarrow \varphi \in X^* : \langle x_n, \varphi \rangle = \varphi(x_n) = 0, \forall n \in \mathbb{N} \Rightarrow \varphi = 0$. The system $\{u_n\}_{n \in \mathbb{N}}$ is minimal in X if and only if there exists a biorthogonal system i.e. there exists a system $\{\vartheta_n\}_{n \in \mathbb{N}} \subset X^*$ such that $\langle u_n, \vartheta_k \rangle = \vartheta_k(u_n) = \delta_{nk}$, where δ_{nk} is the Kronecker symbol.

Basis criterion. The system $\{u_n\}_{n \in \mathbb{N}} \subset X$ is a basis of the space X if and only if the following conditions are satisfied:

- 1) $\{u_n\}_{n \in \mathbb{N}}$ is complete and minimal in X ;
- 2) The projectors

$$P_n f = \sum_{k=1}^n \langle f, \vartheta_k \rangle u_k,$$

where $\{\vartheta_k\}_{k \in \mathbb{N}}$ is a biorthogonal system, are uniformly bounded.

Definition 2.2. A system $\{u_n\}_{n \in \mathbb{N}} \subset X$ is called a basis with brackets in X if there exists a sequence of integers $0 = n_0 < n_1 < n_2 < \dots$, such that each element of $f \in X$ is uniquely represented as a series

$$f = \sum_{k=0}^{\infty} \sum_{i=n_k+1}^{n_{k+1}} c_i u_i,$$

convergent in the norm X .

Basis criterion with brackets. The system $\{u_n\}_{n \in \mathbb{N}} \subset X$ is a basis with brackets of the space X if and only if the following conditions are satisfied:

- 1) $\{u_n\}_{n \in \mathbb{N}}$ is complete and minimal in X ;
- 2) The projectors

$$P_{n_k} f = \sum_{i=1}^{n_k} \langle f, \vartheta_i \rangle u_i,$$

where $\{\vartheta_n\}_{n \in \mathbb{N}}$ is a biorthogonal system, are uniformly bounded.

Proposition 2.1. ([27]) Let the system $\{u_n\}_{n \in \mathbb{N}}$ form a basis with brackets in X and be uniformly minimal, and let the sequence $\{n_{k+1} - n_k\}_{k \in \mathbb{N}}$ be bounded. Then the system $\{u_n\}_{n \in \mathbb{N}}$ forms an usually basis in X .

2.5. Basicity of root functions.

Applying the Fourier method to solution of the problem (B) (or problem (3)-(4)) leads to the following spectral problem

$$\begin{cases} u''(x) + \lambda u(x) = 0, x \in I, \\ u(0) = u(2\pi), u'(0) = 0. \end{cases} \quad (5)$$

The eigenvalues of problem (5) are $\lambda_n = n^2$, $n \in \mathbb{Z}_+$, and corresponding eigenfunctions are $u_{2n}(x) = \cos nx$, $n \in \mathbb{Z}_+$. Each eigenfunction $u_{2n}(x)$, $n \in \mathbb{N}$, has one associated function $u_{2n-1}(x) = x \sin nx$, $n \in \mathbb{N}$. Consider the collection of root functions

$$u_0(x) = 1, u_{2n}(x) = \cos nx, u_{2n-1}(x) = x \sin nx, n \in \mathbb{N}, \quad (6)$$

and also set

$$\begin{aligned} \vartheta_0(x) &= \frac{1}{2\pi^2} (2\pi - x), \vartheta_{2n}(x) = \frac{1}{\pi^2} (2\pi - x) \cos nx, \\ \vartheta_{2n-1}(x) &= \frac{1}{\pi^2} \sin nx, n \in \mathbb{N}. \end{aligned} \quad (7)$$

Note that these systems are biorthogonal conjugate, which can be verified directly. To obtain our main result, we will significantly use the following theorem.

Theorem 2.3 *Let $v \in A_p(I)$, $1 < p < +\infty$. Then the system (6) forms a basis for $L_{p,v}(I)$.*

Proof. To prove the theorem, we first apply the method of work [34]. Estimating the projections $P_{2n}, n \in Z_+$, in a weighted space reduces to estimating the integrals of the Green's function over an expanding system of contours $\Gamma_n = \left\{ \rho \in \mathbb{C}: |\rho| = n + \frac{1}{2}, -\frac{\pi}{2} \leq \arg \rho \leq \frac{\pi}{2} \right\}$ in the complex ρ – plane, which in turn leads to estimating integrals of the Hilbert transform type in a weighted Lebesgue space. Using the boundedness of the Hilbert transform in a weighted Lebesgue space with a weight from the Muckenhoupt class [25], we obtain the uniform boundedness of the system of projectors $\{P_{2n}\}_{n \in Z_+}$ which, according to the basis criterion with brackets, means that system (6) is a basis with brackets in $L_{p,v}(I)$. On the other hand, it is easy to see that the norms of the functions from (6), as well as from (7), are uniformly bounded. Then, applying Proposition 2.1, we obtain that system (6) forms an usually basis in $L_{p,v}(I)$.

3. Main Results

In this section, we will study the existence and uniqueness of strong solution of the problem (B) in the sense of Definition 2.1. For $f(x, y) \in L_{p,v}(\Pi)$ denote $F_n(y) = \langle f(\cdot, y), \vartheta_n(\cdot) \rangle$, i.e.

$$\left. \begin{aligned} F_0(y) &= \frac{1}{2\pi^2} \int_0^{2\pi} f(x, y)(2\pi - x) dx, \\ F_{2n}(y) &= \frac{1}{\pi^2} \int_0^{2\pi} f(x, y)(2\pi - x) \cos nx dx, \\ F_{2n-1}(y) &= \frac{1}{\pi^2} \int_0^{2\pi} f(x, y) \sin nx dx, n \in N. \end{aligned} \right\} \quad (8)$$

The main result of the work is the following theorem.

Theorem 3.1 *Let the weight function $v(x)$ belong to the class $A_p(I)$, $1 < p < +\infty$, and the function $f(x, y) \in W_{p,v}^{1,0}(\Pi)$ and satisfies the following condition $f(0, y) = f(2\pi, y)$. Then problem (B) has a unique solution in $W_{p,v}^2(\Pi)$ and moreover it is valid the following estimate*

$$\|u\|_{W_{p,v}^2(\Pi)} \leq c \|f\|_{W_{p,v}^{1,0}(\Pi)}$$

where $c > 0$ is a constant independent of $f(x, y)$.

Proof. First of all, we note that the uniqueness of the solution follows from the results of [4]. Therefore, we proceed directly to the proof of the existence of a solution. Suppose $u(x, y) \in W_{p;v}^2(\Pi)$ is a solution of the problem (3)-(4). Consider $U_n(y) = \langle u(\cdot, y), \vartheta_n(\cdot) \rangle$, i.e.

$$\left. \begin{aligned} U_0(y) &= \frac{1}{2\pi^2} \int_0^{2\pi} u(x, y) (2\pi - x) dx, \\ U_{2n}(y) &= \frac{1}{\pi^2} \int_0^{2\pi} u(x, y) (2\pi - x) \cos nx dx, \\ U_{2n-1}(y) &= \frac{1}{\pi^2} \int_0^{2\pi} u(x, y) \sin nx dx, \quad n \in \mathbb{N}. \end{aligned} \right\} \quad (9)$$

It is absolutely clear that for a.e. $x \in I$ the relation

$$u(x, y + \delta) - u(x, y) = \int_y^{y+\delta} \frac{\partial u(x, t)}{\partial t} dt, \quad \forall y \in (0, h), \quad (10)$$

holds. Since $\frac{\partial u}{\partial y} \in L_1(\Pi)$, applying Theorem 1.1.1 of [35, p.13], from (10) we obtain that the functions $U_n(y)$ are twice differentiable and can be differentiated under the integral sign. Since the function $u(x, y)$ satisfies the equation (3), multiplying it by $\sin nx$ (by $(2\pi - x) \cos nx$) and integrating over I , we obtain the following relations for $U_{2n-1}(y)$ (respectively, for $U_{2n}(y)$):

$$\left. \begin{aligned} U_0''(y) &= F_0(y), \quad y \in (0, h) \\ U_{2n-1}''(y) - n^2 U_{2n-1}(y) &= F_{2n-1}(y), \quad y \in (0, h), \\ U_{2n}''(y) - n^2 U_{2n}(y) &= -2n U_{2n-1}(y) + F_{2n}(y), \quad y \in (0, h). \end{aligned} \right\} \quad (11)$$

where the functions $F_k(y)$, $k \in \mathbb{Z}_+$ are defined by formulas (8).

Applying the Newton-Leibniz formula and taking into the account the first boundary condition (4), we obtain

$$u_\xi(x) := u(x, \xi) = u(x, 0) + \int_0^\xi \frac{\partial u(x, y)}{\partial y} dy = \int_0^\xi \frac{\partial u(x, y)}{\partial \xi} d\xi, \quad a. e. \quad x \in I.$$

Consequently,

$$|u(x, \xi)| \leq \int_0^\xi \left| \frac{\partial u(x, y)}{\partial y} \right| dy, \quad a. e. \quad x \in I.$$

Hence it immediately follows that

$$\|u_\xi\|_{L_1(I)} = \int_I |u(x, \xi)| dx \leq \int_I \int_0^\xi \left| \frac{\partial u(x, y)}{\partial y} \right| dy dx. \quad (12)$$

We have $|\Pi_\xi| \rightarrow 0$ as $\xi \rightarrow +0$. Then from (12) it follows that

$$u_{\xi}(\cdot) \rightarrow 0, \quad \xi \rightarrow +0, \tag{13}$$

in the norm of the space $L_1(I)$.

Similarly, taking into account the second boundary condition (4), we have

$$u(x, \xi) = u(x, h) - \int_{\xi}^h \frac{\partial u(x, y)}{\partial y} dy = - \int_{\xi}^h \frac{\partial u(x, y)}{\partial y} dy, \quad a. e. \quad x \in I.$$

Hence,

$$\int_I |u(x, \xi)| dx \leq \int_I \int_{\xi}^h \left| \frac{\partial u(x, y)}{\partial y} \right| dy dx. \tag{14}$$

As $|\Pi \setminus \Pi_{\xi}| \rightarrow 0$ when $\xi \rightarrow h - 0$, from (14) it follows that

$$u_{\xi}(\cdot) \rightarrow 0, \quad \xi \rightarrow h - 0, \tag{15}$$

in the norm of the space $L_1(I)$.

On the other hand, it is clear that $U_n(y) \in W_1^2(0, h)$. Hence it immediately follows that there exist the limits

$$\lim_{y \rightarrow +0} U_n(y) = U_n(0), \quad \lim_{y \rightarrow h-0} U_n(y) = U_n(h), \quad \forall n \in Z^+.$$

By (13) and (15), from the last two relations it immediately follows that

$$U_n(0) = 0, \quad U_n(h) = 0, \quad \forall n \in Z_+, \tag{16}$$

Solution of the problem (11), (16) in the case of odd indices is

$$U_{2n-1}(y) = -\frac{1}{n} \frac{\sinh n(h-y)}{\sinh nh} \int_0^y \sinh nt F_{2n-1}(t) dt - \frac{1}{n} \frac{\sinh ny}{\sinh nh} \int_y^h \sinh n(h-t) F_{2n-1}(t) dt, \quad \forall n \in N, \tag{17}$$

and the solution of the problem (11), (16) in the case of even indices is

$$U_0(y) = -\frac{y}{h} \int_0^h (h-t) F_0(t) dt + \int_0^y (y-t) F_0(t) dt, \tag{18}$$

$$U_{2n}(y) = -\frac{1}{n} \frac{\sinh n(h-y)}{\sinh nh} \int_0^y F_{2n}(t) \sinh nt dt - \frac{1}{n} \frac{\sinh ny}{\sinh nh} \int_0^y F_{2n}(t) \sinh nt dt - \frac{2}{n} \frac{\sinh n(h-y)}{\sinh nh} \int_0^y \frac{\sinh n(h-t)}{\sinh nh} \int_0^t F_{2n-1}(\tau) \sinh n\tau d\tau dt -$$

$$\begin{aligned}
 & - \frac{2 \sinh n(h-y)}{n \sinh nh} \int_0^y \frac{\sinh nt}{\sinh nh} \int_t^h F_{2n-1}(\tau) \sinh n(h-\tau) d\tau dt - \\
 & - \frac{2 \sinh ny}{n \sinh nh} \int_y^h \frac{\sinh n(h-t)}{\sinh nh} \int_0^t F_{2n-1}(\tau) \sinh n\tau d\tau dt - \\
 & - \frac{2 \sinh ny}{n \sinh nh} \int_y^h \frac{\sinh nt}{\sinh nh} \int_t^h F_{2n-1}(\tau) \sinh n(h-\tau) d\tau dt, \quad \forall n \in N. \quad (19)
 \end{aligned}$$

Consider the function

$$\begin{aligned}
 u(x, y) = U_0(y) + \sum_{n=1}^{\infty} U_n(y) u_n(x) = U_0(y) + \\
 + \sum_{k=1}^{\infty} (U_{2k}(y) \cos kx + U_{2k-1}(y) x \sin kx), \quad (x, y) \in \Pi, \quad (20)
 \end{aligned}$$

where the coefficients $U_0(y)$, $U_{2k}(\cdot)$, $U_{2k-1}(\cdot)$, $k \in N$, are defined by (17)-(19). Let's show that the function $u(x, y)$ belongs to $W_{p,\nu}^2(\Pi)$. Denote by $u_{\alpha_1, \alpha_2}(x, y)$ the sum of the series obtained by the formal differentiation of the series (20), i.e.

$$u_{\alpha_1, \alpha_2}(x, y) = U_0^{(\alpha_2)}(y) + \sum_{n=1}^{\infty} U_n^{(\alpha_2)}(y) u_n^{(\alpha_1)}(x), \quad (21)$$

where $\alpha_1, \alpha_2 \in Z^+$, $\alpha_1 + \alpha_2 = 0, 1, 2$; $u_{0,0}(x, y) = u(x, y)$.

Let's first consider the series

$$u_1(x, y) = \sum_{k=1}^{\infty} U_{2k-1}(y) x \sin kx.$$

Differentiating this series formally term-by-term, we have

$$\begin{aligned}
 \frac{\partial^2 u_1}{\partial y^2} &= \sum_{k=1}^{\infty} U''_{2k-1}(y) x \sin kx = \sum_{k=1}^{\infty} k^2 U_{2k-1}(y) x \sin kx, \\
 \frac{\partial u_1}{\partial x} &= \sum_{k=1}^{\infty} U_{2k-1}(y) \sin kx + \sum_{k=1}^{\infty} k U_{2k-1}(y) x \cos kx, \\
 \frac{\partial^2 u_1}{\partial x^2} &= 2 \sum_{k=1}^{\infty} k U_{2k-1}(y) \cos kx - \sum_{k=1}^{\infty} k^2 U_{2k-1}(y) x \sin kx. \quad (24)
 \end{aligned}$$

Denote

$$w(x, y) = \sum_{k=1}^{\infty} k^2 U_{2k-1}(y) x \sin kx.$$

Let's show that the function $w(x, y)$ belongs to $L_{p,\nu}(\Pi)$.

From (7), integrating by parts, we obtain

$$F_{2k-1} = \frac{1}{\pi^2} \int_0^{2\pi} f(x; y) \sin kx \, dx = -\frac{1}{\pi^2 k} (f(2\pi; y) - f(0; y)) + \\ + \frac{1}{\pi^2 k} \int_0^{2\pi} f'_x(x; y) \cos kx \, dx = \frac{1}{\pi^2 k} \int_0^{2\pi} f'_x(x; y) \cos kx \, dx = \frac{1}{k} F'_{2k}(y)$$

where is denoted by

$$F'_{2k}(y) = \frac{1}{\pi^2} \int_0^{2\pi} f'_x(x; y) \cos kx \, dx.$$

It is known that if $v \in A_p(I)$, then $\exists \alpha > 1: v \in L_\alpha(I)$ (see, e.g., [36, p. 395]). Let $\frac{1}{\alpha} + \frac{1}{\alpha'} = 1$. Applying Holder's inequality, we obtain

$$\left(\int_0^{2\pi} |w(x, y)|^p v(x) \, dx \right)^{\frac{1}{p}} \leq \left(\int_0^{2\pi} v^\alpha(x) \, dx \right)^{\frac{1}{\alpha}} \left(\int_0^{2\pi} |w(x, y)|^{p\alpha'} \, dx \right)^{\frac{1}{p\alpha'}} \\ = c \left(\int_0^{2\pi} |w(x, y)|^{p_1} \, dx \right)^{\frac{1}{p_1}},$$

where $c = \left(\int_0^{2\pi} v^\alpha(x) \, dx \right)^{\frac{1}{\alpha}}$ (consequently, does not depend on $w(x, y)$) and $p_1 = p\alpha'$. Let us consider separately the cases $p \geq 2$ and $1 < p < 2$.

Let first $p \geq 2$. Then $p_1 = p\alpha' > 2$. From the previous inequality we have

$$\left(\int_0^{2\pi} |w(x, y)|^p v(x) \, dx \right)^{\frac{1}{p}} \leq c \sum_{k=1}^{\infty} |U_{2k-1}(y)| \left(\int_0^{2\pi} |u_{2k-1}(x)|^{p_1} \, dx \right)^{\frac{1}{p_1}} \leq \\ \leq c_1 \sum_{k=1}^{\infty} k^2 |U_{2k-1}(y)| \leq \\ \leq c_1 \sum_{k=1}^{\infty} k^2 \left| -\frac{1}{k} \frac{\sinh k(h-y)}{\sinh kh} \int_0^y \sinh kt F_{2k-1}(t) \, dt \right. \\ \left. - \frac{1}{k} \frac{\sinh ky}{\sinh kh} \int_y^h \sinh k(h-t) F_{2k-1}(t) \, dt \right| \leq$$

$$\begin{aligned} &\leq c_1 \sum_{k=1}^{\infty} \left(\frac{k}{\sinh kh} \left| \left(\sinh k(h-y) \int_0^y \sinh kt F_{2k-1}(t) dt \right) dy \right. \right. \\ &\quad \left. \left. + \left(\sinh ky \int_y^h \sinh k(h-t) F_{2k-1}(t) dt \right) dy \right| \right) \leq \\ &\leq \sum_{k=1}^{\infty} \frac{k}{\sinh kh} \left(\left(\sinh k(h-y) \int_0^y \sinh kt |F_{2k-1}(t)| dt \right) dy \right) \\ &\quad + \left(\sinh ky \int_y^h \sinh k(h-t) |F_{2k-1}(t)| dt \right) dy \end{aligned}$$

Hence, first integrating with respect to $y \in (0, h)$ and then applying Holder's inequality for any $\beta \in (1, \infty)$, we obtain

$$\begin{aligned} \|w\|_{L_{p,\nu}(\Pi)} &\leq c_1 \sum_{k=1}^{\infty} \left(\frac{k}{\sinh kh} \left(\int_0^h (\sinh k(h-y) \int_0^y \sinh kt |F_{2k-1}(t)| dt) dy \right) \right. \\ &\quad \left. + \left(\int_0^h \sinh ky \int_y^h \sinh k(h-t) |F_{2k-1}(t)| dt \right) dy \right) \leq \\ &\leq c_1 \sum_{k=1}^{\infty} \int_0^h |F_{2k-1}(t)| dt \leq c_1 \int_0^h \left(\sum_{k=1}^{\infty} |F_{2k-1}(t)| \right) dt \leq \\ &\leq c_2 \int_0^h \left(\sum_{k=1}^{\infty} \frac{1}{k} |F'_{2k}(t)| \right) dt \leq c_2 \left(\sum_{k=1}^{\infty} \frac{1}{k^{\beta'}} \right)^{\frac{1}{\beta'}} \int_0^h \left(\sum_{k=1}^{\infty} |F'_{2k}(t)|^{\beta} \right)^{\frac{1}{\beta}} dt \end{aligned}$$

Now, assuming $\beta \geq 2$ and applying classical Hausdorff-Young inequality (see, e.g. [37, p.154]), we obtain

$$\|w\|_{L_{p,\nu}(\Pi)} \leq c_3 \int_0^h \left(\left\| \frac{\partial f}{\partial x} \right\|_{L_{\beta'}(I)} \right) dt. \tag{25}$$

It is known (see, e.g. [25,26]) that if $\nu \in A_p(I)$, $1 < p < +\infty$, then $\exists q: 1 < q < p \Rightarrow \nu \in A_q(I)$. Let $r = \frac{p}{q}$ and $g \in L_{p,\nu}(I)$. Then $1 < r < p$ and we have

$$\begin{aligned} \left(\int_I |g|^r dx \right)^{\frac{1}{r}} &= \left(\int_I |g|^{\frac{p}{q}} v^{\frac{1}{q}} v^{-\frac{1}{q}} dx \right)^{\frac{1}{r}} \leq \left(\int_I |g|^p v dx \right)^{\frac{1}{qr}} \left(\int_I v^{-\frac{q'}{q}} dx \right)^{\frac{1}{qr}} = \\ &= \left(\int_I |g|^p v dx \right)^{\frac{1}{p}} \left(\int_I v^{-\frac{1}{q-1}} dx \right)^{\frac{q-1}{p}}. \end{aligned}$$

As $-\frac{q'}{q} = \frac{1}{1-q}$, from $v \in A_q(I)$ it follows that $v^{-\frac{1}{q-1}} \in L_1(I)$. Then, the last inequality means $g \in L_r(I)$ and

$$\|g\|_{L_r(I)} \leq c \|g\|_{L_{p,v}(I)},$$

where $c > 0$ is a constant independent of g . Also note that the continuous embedding $L_{p,v}(I) \subset L_\alpha(I)$ is true for every $\alpha \in (1, r)$. Let's choose β big enough to satisfy the condition $1 < \beta' < r$. Then from (25) we obtain

$$\|w\|_{L_{p,v}(\Pi)} \leq c \left(\left\| \frac{\partial f}{\partial x} \right\|_{L_{p,v}(\Pi)} \right)$$

Now let $p \in (1, 2)$. As in the previous case, note that there exists a number $\alpha > 1$ such that $v \in L_\alpha(I)$. But then $v \in L_s(I)$ for every $s \in (1, \alpha)$. Therefore, choosing $\alpha > 1$ close enough to 1, we can provide that $p_1 = p \alpha' > 2$ (this is possible, because $\alpha' \rightarrow +\infty$ as $\alpha \rightarrow 1 + 0$). With this, further considerations are carried out similar to the previous case.

Other series from (21)-(23), and, consequently, all series from (20) are estimated in a similar way. So, as a result, we obtain

$$\|u\|_{W_{p,v}^2(\Pi)} \leq c \left(\left\| \frac{\partial f}{\partial x} \right\|_{L_{p,v}(\Pi)} \right),$$

where $c > 0$ is a constant independent of f . The fulfillment of equation (3) for $u(\cdot; \cdot)$ is verified directly. The fulfillment of boundary conditions (4) follows from the relations for the trace of this function, proved in our previous work [4].

The theorem is proved.

Thus, based on Theorem 3.1 of this article and our previous work [4], we obtain the validity of the following theorem on the correct solvability of problem (1)-(2).

Theorem 3.2 *Let the weight function $v(x)$ belong to the class $A_p(I)$, $1 <$*

$p < \infty$ and, and the boundary functions $\varphi(x)$ and $\psi(x)$ belong to $W_{p;\nu}^2(I)$ and satisfy the following conditions

$$\varphi(0) - \varphi(2\pi) = \varphi'(0) = \psi(0) - \psi(2\pi) = \psi'(0) = 0,$$

and the function $f(x, y) \in W_{p;\nu}^{1,0}(\Pi)$ and satisfies the condition $f(0, y) = f(2\pi, y)$. Then problem (1)-(2) has a unique solution in $W_{p;\nu}^2(\Pi)$ and moreover it is valid the following estimate

$$\|u\|_{W_{p;\nu}^2(\Pi)} \leq c \left(\|\varphi\|_{W_{p;\nu}^2(I)} + \|\psi\|_{W_{p;\nu}^2(I)} + \|f\|_{W_{p;\nu}^{1,0}(\Pi)} \right),$$

where $c > 0$ is a constant independent of φ, ψ and f .

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